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# Suppression of $T_c$ in the $(Y_{0.9}Ca_{0.1})Ba_2Cu_{4-x}Fe_xO_8$ system

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## Abstract

In this paper, the effects produced by iron substitutions in the  $(Y_{0.9}Ca_{0.1})Ba_2Cu_{4-x}Fe_xO_8$  system on the superconducting and structural properties are studied. The Rietveld-fit of the crystal structure and Mössbauer spectroscopy results of  $(Y_{0.9}Ca_{0.1})Ba_2Cu_{4-x}Fe_xO_8$  samples indicate that the iron atoms occupy the Cu(1) sites of the  $(Cu-O)_2$  double chains in fivefold coordination at low iron concentrations. Besides at high iron concentrations the iron atoms occupy the Cu(1) sites of single Cu–O chains and Cu(2) sites in the  $CuO_2$  planes of the  $(Y_{0.9}Ca_{0.1})Ba_2Cu_4O_8$  phase with structural defects. Simultaneously, as iron concentration increases, a faster decrease of  $T_c$  is observed in this material compared with the  $YBa_2Cu_{3-x}Fe_xO_{7-\delta}$  system. According to the charge-transfer model proposed for  $YBa_2Cu_4O_8$  under pressure, the decrease in the Cu(1)–O(4) bond length in parallel to the increase in the Cu(2)–O(4) bond length may affect the charge-transfer mechanism leading to the suppression of  $T_c$ .

## 1. Introduction

Among the high- $T_c$  superconductors, the  $YBa_2Cu_4O_8$  (Y124,  $T_c = 80$  K) system has been well studied with regard to structural and physical properties. As is now well-known, the Y124 crystal structure is closely related to that of the  $YBa_2Cu_3O_7$  (Y123,  $T_c = 90$  K) with a double chain  $(Cu-O)_2$ , instead of a single one, along the  $b$ -axis of the unit cell. The positions of the Cu atoms in adjacent Cu–O chains differ by a distance  $b/2$  along the  $b$ -axis [1, 2]; this fact leads to a  $c$  length parameter longer than in the Y123 by 27.25 Å. The Y substitution by Ca in the Y124 phase gives rise to the increase of the superconducting transition temperature  $T_c$  to 90 K for the nominal composition of  $(Y_{0.9}Ca_{0.1})Ba_2Cu_4O_8$  [3]. As was suggested previously [3] the doped  $Ca^{2+}$  ions introduce additional holes as charge carriers into the Y124 phase which increase the value of  $T_c$ .

On the other hand, since the discovery of high-temperature superconductors many experiments have been carried out in order to determine the effects of Cu substitution by other 3d metals on the superconductivity. It was hoped that these investigations would reveal clues on the mechanism that causes the high transition temperatures in these materials.

However, it is still not clear what causes the suppression of  $T_c$  as the dopant concentration is increased. Among the most common effects proposed in the literature are magnetic pair breaking, change in the local symmetry and purely electronic dynamical mechanisms [4]. In any case, basic issues such as dopant site localization and charge state must be considered. Mössbauer spectroscopy has been proved to be a powerful tool in determination of the oxidation state of the substituent Fe atoms as well as their environment (i.e. crystallographic site in the structure). Although it is now apparent that the details of the sample as well as of oxygen stoichiometry, a considerable amount of experimental evidence appears to be at least qualitatively well established.

In the present research we have carried out systematic substitutional studies on  $(Y_{0.9}Ca_{0.1})Ba_2Cu_{4-x}Fe_xO_8$  ( $x = 0.025, 0.05, 0.075, 0.1$  and  $0.2$ ) to clarify where the iron atoms are situated in the structure and how they affect its crystalline structure and transport properties. We present x-ray diffractograms, Mössbauer spectroscopy, resistance versus temperature and ac susceptibility versus temperature measurements results. In order to attain a site assignment, the Mössbauer results were complemented with a Rietveld refinement of the x-ray diffraction data.

## 2. Experimental details

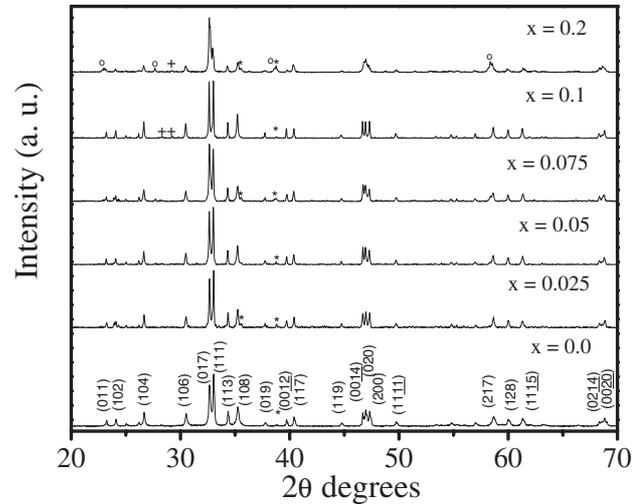
We have synthesized  $(Y_{0.9}Ca_{0.1})Ba_2Cu_{4-x}Fe_xO_8$  (YCa124Fe) samples with  $x = 0, 0.025, 0.05, 0.075, 0.1$  and  $0.2$  with ambient oxygen pressure using a combined method described as follows: (a) powders of  $Y_2O_3$ ,  $CaCO_3$ ,  $BaCO_3$ ,  $CuO$  and  $^{57}Fe_2O_3$  (99.999%) were mixed in stoichiometric amounts, diluted in 10 ml of nitric acid (96%) and heated until all the liquid evaporated; the resulting green paste was fired at  $700^\circ C$  in air to obtain small grain ( $1\text{--}10\ \mu m$ ) precursors. (b) The calcinated powders were mixed with sodium nitrate powder (20% in weight) as catalyst to accelerate the reaction, pressed into pellets, heat treated at  $810^\circ C$  for 20 h in oxygen flow and slowly cooled ( $0.5^\circ C\ min^{-1}$ ) to room temperature. The heat treatment was repeated at least four times.

Phase identification of the samples was done with an x-ray diffractometer Siemens D5000 using  $Cu\text{-}K\alpha$  radiation and a Ni filter. Intensities were measured in steps of  $0.02^\circ$  for 14 s in the  $2\theta$  range  $20^\circ\text{--}70^\circ$  at room temperature and crystallographic phases were identified by comparison with the x-ray patterns of the JCPDS database. Crystallographic parameters were refined using a Rietveld refinement program, RIETQUAN v2.3 [5], with multi-phase capability. The superconducting transition temperatures were determined in a closed-cycle helium refrigerator by measuring resistance versus temperature curves by the standard four-probe technique. The ac magnetization against temperature behaviour in the temperature range of  $12\text{--}300\ K$ , was obtained using an APD Cryogenics Inc. Superconductor Characterization Cryostat in conjunction with an EG&G model 5210 lock-in amplifier. Measurements were taken at a frequency of 125 Hz in an ac field of 1.1 mT. The room temperature Mössbauer spectra were recorded in transmission geometry with a constant acceleration spectrometer, using a  $^{57}Co$  in Rh source. The spectra were fitted with a constrained-least-squares programme.

## 3. Results and discussion

Figure 1 shows x-ray diffraction patterns for the  $(Y_{0.9}Ca_{0.1})Ba_2Cu_{4-x}Fe_xO_8$  (YCa124Fe) samples synthesized. The analysis of these data indicates that all samples correspond to the Y124 structure although for  $0 \leq x \leq 0.075$  (low concentration) faint features of  $CuO$  (IDDE no 05-661) appear and for  $x \geq 0.1$  (high concentration) additional peaks corresponding to  $(Y_{0.9}Ca_{0.1})Ba_2Cu_3O_y$  (YCa123) and to  $BaCuO_2$  (IDDE no 38-1402) impurities can also be observed. Yanagizawa *et al* [6] have shown, by high resolution transmission electron microscopy (HRTEM), that at around a 3% level of iron doping, a destabilization of the  $(Cu\text{-}O)_2$  double chains of Y124 takes place, which produces random distributed  $Cu\text{-}O$  single chains in the Y124 structure. These structural defects are manifest in the x-ray diffraction pattern for the presence of traces associated with ‘Y123’ pseudo-phase. This effect would be difficult to understand if the site occupancy of the iron atoms is other than the  $Cu(1)$  sites. In this context, the Y124 phase with structural defects is a ‘Y123’ pseudo-phase.

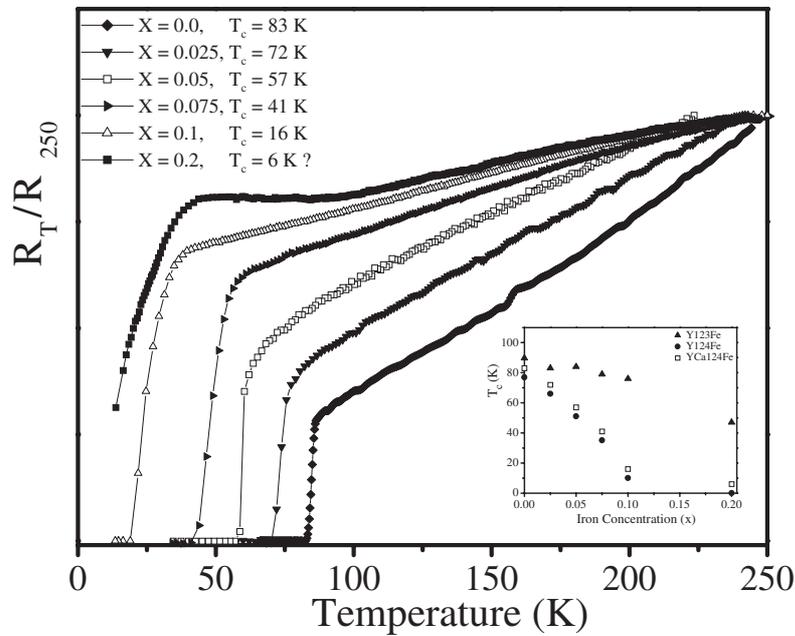
Figure 2 shows the normalized resistance versus temperature of the studied samples. We observed that  $T_c$



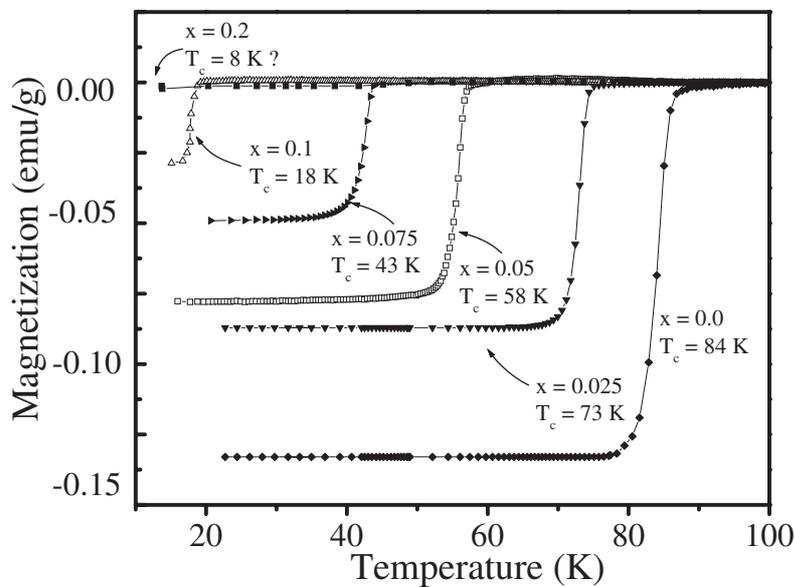
**Figure 1.** X-ray diffraction patterns of the  $(Y_{0.9}Ca_{0.1})Ba_2Cu_{4-x}Fe_xO_8$  (YCa124Fe) samples. The symbols represent: (\*)  $CuO$ , (+)  $BaCuO_2$ , (o) ‘YCa123Fe’.

decreases continuously increasing the iron concentration ( $x$ ). The zero resistance of the sample with  $x = 0.2$ , is an extrapolated value, and it is indicated as a solid square in the inset of figure 2. In this inset, the  $T_c$  variation with  $x$  for this system compared with the  $YBa_2Cu_{4-x}Fe_xO_8$  (Y124Fe) [7] and  $YBa_2Cu_{3-x}Fe_xO_y$  (Y123Fe) [8, 9] systems is shown. We observed that the  $T_c$  values in YCa124Fe system are larger than those reported in [7]. According to Miyatake *et al* [3] the reason of this increase in  $T_c$  is the introduction of holes as charge carriers due to the substitution of Ca ions in the Y sites. On the other hand, we also observed in the Y124Fe and YCa124Fe systems, that at  $x \leq 0.1$ , the rate of decrease of  $T_c$  with iron concentration is linear. But at  $x = 0.2$ , the YCa124Fe system continues being superconductor ( $T_c \sim 6\ K$ ) while in the Y124Fe system  $T_c$  is suppressed. Figure 3 shows magnetization as a function of temperature for the YCa124Fe samples. We observed that the values of the transition temperature determined from the magnetization measurement are in good agreement with the transition temperatures from the resistivity versus temperature curves.

The diffraction patterns of YCa124Fe samples were Rietveld-fitted using an orthorhombic structure with space group  $Ammm$  (no 64). In the refinement process: (a) it was assumed that Ca atoms substitute the Y sites, (b) the presence of different impurities, such as  $BaCuO_2$ ,  $CuO$  phases and ‘YCa123’ pseudo-phase was considered, (c) the occupancy factor of Cu and Fe atoms in the  $Cu(1)$  and  $Cu(2)$  sites was varied independently in order to allow the program to establish the site occupancy of the Fe atoms in the YCa124 and ‘YCa123’ structures, (d) we permitted the possibility of extra oxygen atoms in the structure: O(5), between two adjacent double chains planes in front of a  $Cu(1)$  site, along the  $a$ -axis, (e) finally, the isotropic thermal parameters (in  $\text{Å}^2$ ) were kept fixed at the values obtained by Adachi *et al* [10]. In figure 4(a), the results for the sample with  $x = 0.025$  as representative of the low concentration samples are shown. In similar way, figure 4(b) presents the results for sample with  $x = 0.2$  as representative of the high concentration samples. In tables 1 and 2 the structural parameters and occupation factors obtained



**Figure 2.** Normalized resistance versus temperature of  $(Y_{0.9}Ca_{0.1})Ba_2Cu_{4-x}Fe_xO_8$  (YCa124Fe) samples. The inset shows the  $T_c$  variation with iron concentration  $x$ .

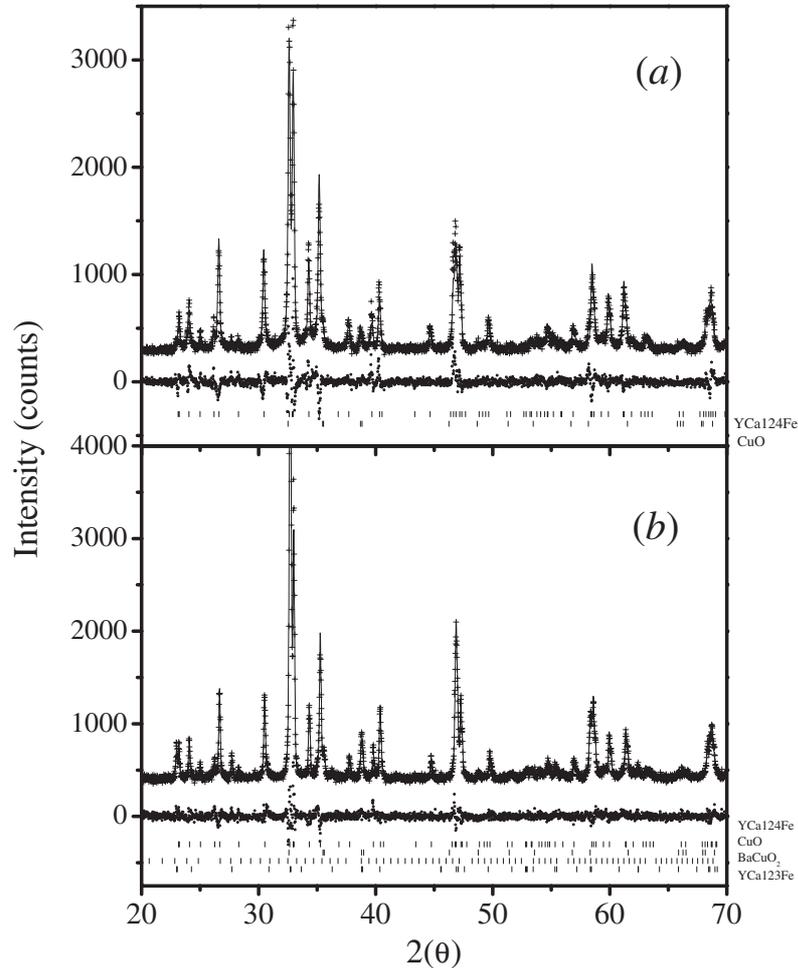


**Figure 3.** Magnetization as a function of temperature for the  $(Y_{0.9}Ca_{0.1})Ba_2Cu_{4-x}Fe_xO_8$  (YCa124Fe) samples.

for  $x = 0.025$  and  $x = 0.2$  are summarized. O(1) refers to the oxygen atoms along the double chains, O(2) and O(3) to the oxygen atoms in the  $CuO_2$  planes and O(4) to the oxygen atoms in the Ba atoms plane. Our results of Rietveld-fit, show that the  $a$ -axis and  $b$ -axis lattice parameters increase slightly, as a function of increasing ( $x$ ) and the  $c$ -axis one decreases continuously (see figure 5). We observed that the structure of the YCa124Fe samples remains orthorhombic up to the highest Fe doping ( $x = 0.2$ ).

On the other hand, the refinement of atomic positions and occupation factor shows that at low iron concentration the iron atoms are fully localized on the Cu(1) site of the double chains of the YCa124Fe structure. All attempts to put some iron on the Cu(2) site lead to a negative value of

this parameter after refinement. At high iron concentrations, our results show that the iron atoms substitute the Cu(1) sites of the YCa124Fe structure and the Cu(1) and Cu(2) sites of the 'YCa123Fe' (see tables 1 and 2). These replacements originate significant changes in the Cu–O bond lengths. Figure 6 shows the Cu(1)–O(4) and Cu(2)–O(4) bond lengths obtained from the refined atomic positions. We observed that the former decreases whilst the latter increases with iron concentration. Moreover, with increasing iron concentration, the occupation number for O(5) oxygen of the YCa124Fe structure is increased in good agreement with previous studies in the  $YBa_2Cu_{3-x}Co_xO_{7-y}$  system [11]. In this case, for low cobalt concentration, the Co atoms preferentially occupy the Cu(1) site and simultaneously attracting extra



**Figure 4.** (a) Rietveld refinement on the x-ray diffraction pattern for the  $x = 0.025$  sample. (b) Rietveld refinement on the x-ray diffraction pattern for the  $x = 0.2$  sample. Experimental pattern (dots), calculated pattern (continuous line), their difference (middle line) and the calculated peaks positions (bottom).

oxygen into these layers in an amount consistent with  $Co^{3+}$ . As a consequence of this replacement, the  $Cu(1)-O(4)$  and  $Cu(2)-O(4)$  bond lengths change in the same way.

To supplement the site assignment of iron atoms in the YCa124Fe structure, we have performed studies of Mössbauer spectroscopy. Figures 7(a) and (b) show the Mössbauer spectra for  $x = 0.025$  and  $x = 0.2$ , respectively. Each of the spectra for the three lower iron concentrations consists of a single quadrupole doublet-labelled A in figure 7(a)—whose parameters (quadrupole splitting  $\Delta Q$ , isomer shift (IS) and line width  $\Gamma$ ) remain essentially constant (see table 3). However, the spectra of the two higher iron concentration samples consist of four quadrupole doublets-labelled A, B, C and D (see figure 7(b), the most prominent of which (A) corresponds to the single doublet observed at low iron concentrations. Our study shows that at low and high iron concentrations, the observed Mössbauer spectra have  $\Delta Q$  values close to those observed in the Y124Fe system (see table 3) [7]. These results suggest that the  $\Delta Q$  values are independent of the substitution of Ca atoms (10%) in the Y sites. Based on the above mentioned, we propose the following assignment: at low iron concentrations, the  $0.79 \text{ mm s}^{-1}$  quadrupole doublet (A) value is associated with  $Fe^{3+}$  atoms in the Cu(1) sites in fivefold

coordination, resulting from extra oxygen atoms, named O(5), that are attracted by Fe atoms at the Cu(1) sites, and occupy sites placed along the  $a$ -axis between two  $(Cu-O)_2$  double chains (see tables 1 and 2), similar site assignment of the iron atoms has been shown by Felner *et al* [12, 13].

As is now well known, the Mössbauer spectra reported in the literature for the Y123Fe phase consist of three main quadrupole doublets, which have  $\Delta Q$  of approximately  $2.0 \text{ mm s}^{-1}$ ,  $1.10 \text{ mm s}^{-1}$  and  $0.55 \text{ mm s}^{-1}$ . The only doublet that has been associated with the Cu(2) sites is the one whose  $\Delta Q$  is around  $0.55 \text{ mm s}^{-1}$  [14]. At high iron concentrations, besides the presence of quadrupole doublet (A), the quadrupole splittings of doublets B ( $2.0 \text{ mm s}^{-1}$ ) and C ( $0.90 \text{ mm s}^{-1}$ ) have values close to those associated with  $Fe^{3+}$  occupying Cu(1) of the Cu–O single chains of the Y123 structure that result from the destabilization of the  $(Cu-O)_2$  double chains [9]. The doublet (B) is associated with  $Fe^{3+}$  in the Cu(1) sites with fourfold (planar) coordination and the doublet (C), associated with  $Fe^{3+}$  also in the Cu(1) sites with fourfold (non-planar) coordination and/or with fivefold (pyramidal) coordination. While the doublet (D) with  $\Delta Q$  of around  $0.56 \text{ mm s}^{-1}$  is associated with the Cu(2) sites.

In order to find a more in-depth correlation between the structure and  $T_c$  of the YCa124Fe system, the evolution of  $T_c$

**Table 1.** Rietveld refinement results for the  $(Y_{0.9}Ca_{0.1})Ba_2Cu_{4-x}Fe_xO_8$  (YCa124Fe) structure with (a)  $x = 0.025$  and (b)  $x = 0.2$ .

Atom	X	Y	Z	N
(a) $x = 0.025^a$				
Cu(1)	0	0	0.2127(3)	1.880(9)
Fe(1)	0	0	0.2127(3)	0.051(7)
Cu(2)	0	0	0.0639(3)	1.965(7)
Ba	0.5	0.5	0.1344(3)	2.0
Y	0.5	0.5	0	0.977(2)
Ca	0.5	0.5	0	0.061(2)
O(1)	0	0.5	0.2164(1)	2.0
O(2)	0.5	0	0.0590(1)	2.0
O(3)	0	0.5	0.0520(2)	1.947(3)
O(4)	0	0	0.1400(1)	2.0
(b) $x = 0.2^b$				
Cu(1)	0	0	0.2127(3)	1.880(9)
Fe(1)	0	0	0.2127(3)	0.051(7)
Cu(2)	0	0	0.0639(3)	1.965(7)
Ba	0.5	0.5	0.1345(2)	2.0
Y	0.5	0.5	0	0.977(2)
Ca	0.5	0.5	0	0.061(2)
O(1)	0	0.5	0.2168(2)	2.0
O(2)	0.5	0	0.0590(1)	2.0
O(3)	0	0.5	0.0520(2)	1.947(3)
O(4)	0	0	0.1414(2)	2.0
O(5)	0.5	0	0.2106(3)	0.024(4)

<sup>a</sup> The calculated crystal parameters (in Å) are:  $a = 3.8461(2)$ ,  $b = 3.8724(2)$  and  $c = 27.256(1)$ . Isotropic thermal parameters (in Å<sup>2</sup>): 0.005 for cations, 0.010 for oxygen (see [10]). The result reliability factors were:  $R_{wp} = 10\%$ ,  $R_E = 6.2\%$ ,  $R_b = 11.4\%$ ,  $\chi^2 = 1.6\%$ .

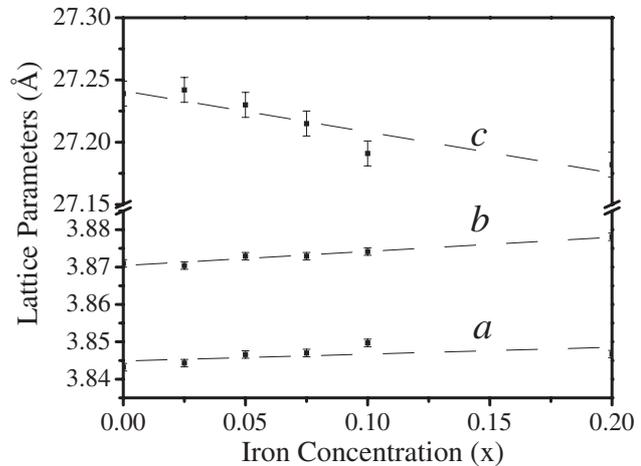
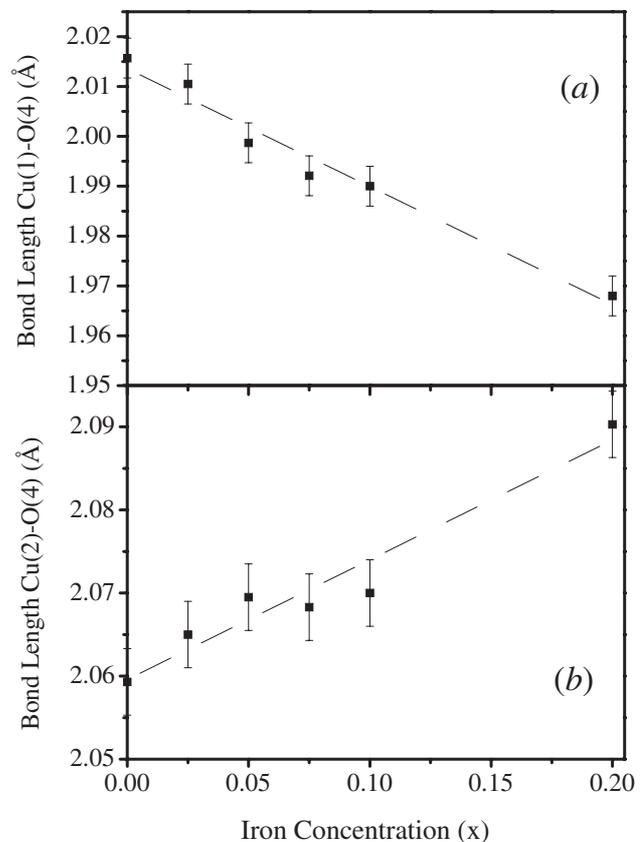
<sup>b</sup> The calculated crystal parameters (in Å) are:  $a = 3.8472(3)$ ,  $b = 3.8789(2)$  and  $c = 27.205(2)$ . Isotropic thermal parameters (in Å<sup>2</sup>): 0.005 for cations, 0.010 for oxygen (see [10]). The result reliability factors were:  $R_{wp} = 5.8\%$ ,  $R_E = 4.3\%$ ,  $R_b = 5.9\%$ ,  $\chi^2 = 1.34\%$ .

**Table 2.** Rietveld refinement results for the YCa123Fe impurity structure with  $x = 0.2^a$ .

Atom	X	Y	Z	N
Cu(1)	0	0	0	0.979(2)
Fe(1)	0	0	0	0.074(2)
Cu(2)	0	0	0.3553(7)	1.976(4)
Fe(2)	0	0	0.3553(7)	0.002(3)
Ba	0.5	0.5	0.1840(3)	2.0
Y	0.5	0.5	0.5	0.928(3)
Ca	0.5	0.5	0.5	0.026(2)
O(1)	0	0.5	0	1.775(4)
O(2)	0.5	0	0.3686(2)	2.0
O(4)	0	0	0.1629(2)	2.0

<sup>a</sup> The space group is  $P4/mmm$  and the calculated crystal parameters (in Å) are:  $a = b = 3.8775(9)$  and  $c = 11.6016(6)$ . The O(2) and O(3) positions and occupation factors are equivalent in the tetragonal structure.

and Cu(1)–O(4) bond length as a function of iron concentration ( $x$ ) is displayed in figure 8. We observed that the Cu(1)–O(4) bond length decreases as the iron concentration is increased. The change in bond lengths and  $T_c$  with iron concentration,

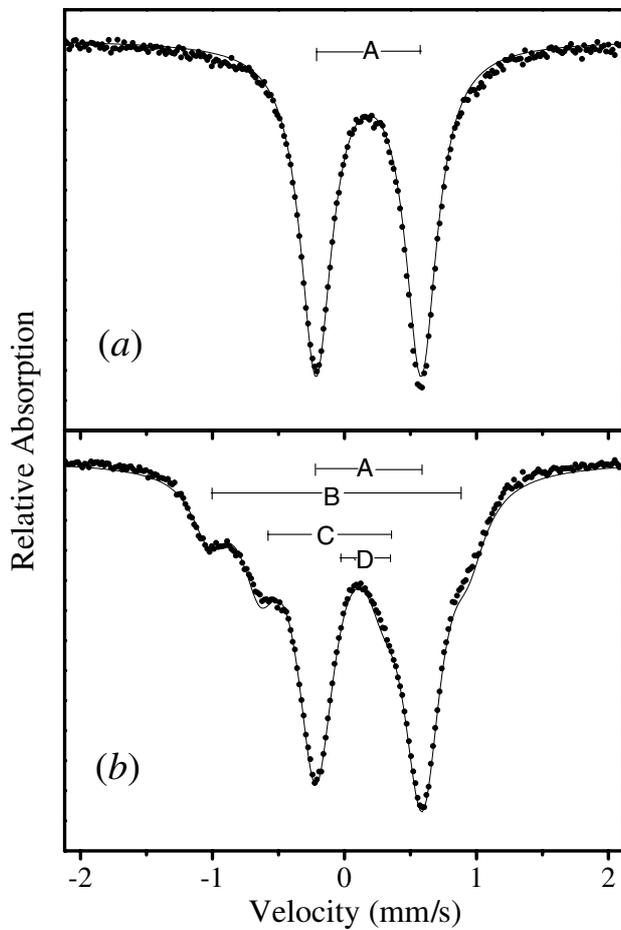
**Figure 5.** Crystal lattice parameters as a function of iron concentration  $x$ .**Figure 6.** Cu(1)–O(4) and Cu(2)–O(4) bond lengths versus iron concentration  $x$ .

suggests that significant charge-transfer effects occur. It has been claimed that the enhancement of  $T_c$  in Y124 under pressure is mainly due to the shortening of the Cu(2)–O(4) bond length and relative elongation of the Cu(1)–O(4) bond length. These changes in the bond lengths were related to the charge-carriers transfer (using the bond valence sum formalism [15, 16]) from the Cu–O chains to the CuO<sub>2</sub> plane, which causes an increase in charge carriers in the conducting plane CuO<sub>2</sub> leading to the increase in  $T_c$  [17–19].

**Table 3.** Mössbauer parameters of the  $(Y_{0.9}Ca_{0.1})Ba_2Cu_{4-x}Fe_xO_8$  (YCa124Fe) samples<sup>a</sup>.

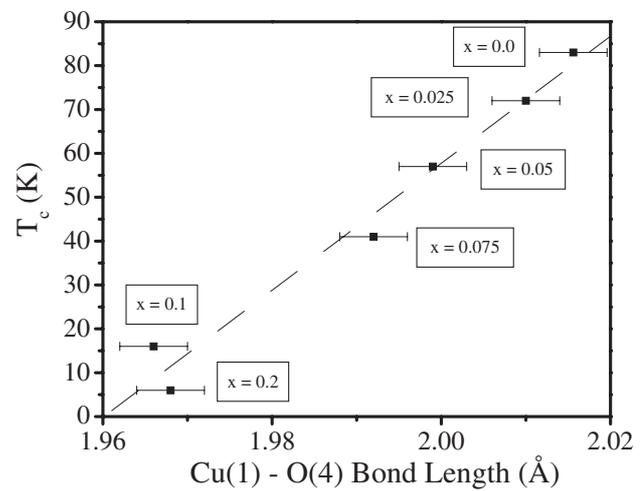
$X$	$IS_A$	$\Delta Q_A$	$\Gamma_A$	$IS_B$	$\Delta Q_B$	$\Gamma_B$	$IS_C$	$\Delta Q_C$	$\Gamma_C$	$IS_D$	$\Delta Q_D$	$\Gamma_D$
$Y_{0.9}Ca_{0.1}Ba_2Cu_{4-x}Fe_xO_8$ structure												
0.025	0.32	0.796	0.287									
0.050	0.32	0.78	0.28									
0.075	0.32	0.79	0.29									
0.1	0.30	0.77	0.26	0.07	1.98	0.35	0.002	0.99	0.31	0.02	0.50	0.31
0.2	0.306	0.802	0.28	-0.05	1.95	0.349	-0.073	1.131	0.35	-0.021	0.596	0.35
$YBa_2Cu_{4-x}Fe_xO_8$ structure [7]												
0.025	0.32	0.78	0.28									
0.050	0.32	0.78	0.28									
0.075	0.32	0.79	0.29									
0.1	0.30	0.77	0.26	0.07	1.98	0.35	0.002	0.99	0.31	0.02	0.50	0.31
0.2	0.30	0.79	0.31	0.07	2.02	0.33	-0.06	0.91	0.32	-0.07	0.56	0.27

<sup>a</sup> The isomer shifts are with respect to iron and all the figures are in  $mm\ s^{-1}$ .



**Figure 7.** Mössbauer spectra (a) for  $x = 0.025$  and (b)  $x = 0.2$ .

Comparing our results with Y124 system under pressure, we observed that the substitution of Fe atoms in Cu(1) sites, induces an increase in the Cu(2)–O(4) bond length and a decrease in the Cu(1)–O(4) bond length. Simultaneously  $T_c$  decreases with increase in iron concentration. According to the charge-transfer model proposed for Y124 under pressure, this behaviour in the Cu(1)–O(4) and Cu(2)–O(4) bond lengths of YCa124Fe system may induce an opposite charge transfer, i.e. charge transfer from the  $CuO_2$  plane to the double chains



**Figure 8.**  $T_c$  versus Cu(1)–O(4) bond length for each concentration of iron.

(Cu–O)<sub>2</sub>. This opposite charge transfer can reduce the charge carriers in the conducting  $CuO_2$  plane leading to the suppression of  $T_c$ . This may be one of the reasons for the  $T_c$  decrease in YCa124Fe.

#### 4. Conclusions

We have presented a detailed crystallographic, transport and magnetic characterization of the  $(Y_{0.9}Ca_{0.1})Ba_2Cu_{4-x}Fe_xO_8$  (YCa124Fe) system by means of XRD powder at room temperature, Mössbauer spectroscopy, resistance versus temperature curves and magnetization measurements. Our results of Mössbauer spectroscopy and Rietveld-fit, indicate that at low iron concentrations the iron atoms occupy partially the Cu(1) sites of the double (Cu–O)<sub>2</sub> chains and at high iron concentrations also occupy partially the Cu(1) and Cu(2) sites of the YCa124Fe with structural defects ('YCa123Fe'). These replacements induce significant changes both in the Cu(1)–O(4) and Cu(2)–O(4) bond lengths, the former decreases whilst the latter increases with iron concentration. Simultaneously our resistance versus temperature and ac susceptibility magnetic measurements show that the transition temperatures of the YCa124Fe samples decrease with the

increasing iron concentration  $x$ . The change in bond lengths and  $T_c$  with the iron concentration, suggest that significant charge-transfer effects are occurring. According to the charge-transfer model proposed for Y124 under pressure, one should expect a less effective charge transfer in the YCa124Fe system, leading to the suppression of  $T_c$ . In summary, these facts suggest that the substitution of Fe atoms in the Cu(1) sites modifies the charge-transfer mechanism through changes in the Cu(2)–O(4) and Cu(1)–O(4) bond lengths.

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### References

- [1] Zandbergen H, Gronsky R, Wang K and Thomas G 1988 *Nature* **331** 596
- [2] Marshall A F, Barton R W, Char K, Kapitulnik A, Oh B, Hammond R H and Laderman S S 1988 *Phys. Rev. B* **37** 9353
- [3] Miyatake T, Gotoh S, Koshizuka N and Tanaka S 1989 *Nature* **341** 41
- [4] Smith M G, Taylor R D and Thompson J D 1993 *Physica C* **208** 91
- [5] Lutterotti L, Scardi P and Maistrelli P 1992 *J. Appl. Cryst.* **25** 459
- [6] Yanagisawa K, Matsui Y, Kodama Y, Yamada Y and Matsumoto T 1991 *Physica C* **183** 197
- [7] Yanagisawa K, Matsui Y, Kodama Y, Yamada Y and Matsumoto T 1992 *Physica C* **191** 32
- [8] Akachi T, Escamilla R, Marquina V, Jiménez M, Marquina M L, Gómez R, Ridaura R and Aburto S 1998 *Physica C* **301** 315
- [9] Tarascon J M, Barboux P, Miceli P F, Greene L H, Hull G W, Eibschutz M and Sunshine S A 1988 *Phys. Rev. B* **37** 7458
- [10] Maeno Y, Kato M, Auki Y and Fujita T 1987 *Japan. J. Appl. Phys.* **26** L1982
- [11] Adachi S, Watanabe N, Seiji N, Koshizuka N and Yamauchi H 1993 *Physica C* **207** 127
- [12] Miceli P F, Tarascon J M, Greene L H, Barboux P, Rotella F J and Jorgensen J D 1988 *Phys. Rev. B* **37** 5932
- [13] Felner I and Brosh B 1991 *Phys. Rev. B* **43** 10364
- [14] Felner I, Novik I, Brosh B, Hechel D and Bauminger E R 1991 *Phys. Rev. B* **43** 8737
- [15] Dunlap B D, Jorgensen J D, Segre C, Dwight A E, Matykiewicz J L, Lee H, Peng W and Kimball C W 1989 *Physica C* **158** 397
- [16] Brown I D 1989 *J. Solid State Chem.* **82** 122
- [17] Jorgensen J D, Pei S Y, Lightfoot P, Hinks D G, Veal B W, Dabrowski B, Paulikas A P and Kleb R 1990 *Physica C* **171** 93
- [18] Kaldis E, Fischer P, Hewat A W, Hewat E A, Karpinski J and Rusiecki S 1989 *Physica C* **159** 668
- [19] Nelmes R J, Loveday J S, Kaldis E and Karpinski J 1992 *Physica C* **172** 311
- [20] Yamada Y, Jorgensen J D, Pei S Y, Lightfoot P, Kodama Y, Matsumoto T and Izumi F 1991 *Physica C* **173** 185